Как «увидеть» скрытую квантовую критическую точку?

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Quantum criticality in Mn_{1-x}Fe_xSi. Theoretical expectations.



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concentration, etc.)



PHYSICAL REVIEW B 82, 064404 (2010)

Quantum phase transitions in single-crystal $Mn_{1-x}Fe_xSi$ and $Mn_{1-x}Co_xSi$: Crystal growth, magnetization, ac susceptibility, and specific heat PHYSICAL REVIEW B 83, 224411 (2011)

Chiral criticality in the doped helimagnets Mn_{1-y}Fe_ySi

Sergey V. Grigoriev,¹ Evgeny V. Moskvin,¹ Vadim A. Dyadkin,¹ Daniel Lamago,^{2,3} Thomas Wolf,³ Helmut Eckerlebe,⁴ and Sergey V. Maleyev¹





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Correlation between magnetic susceptibility and polarized neutron scattering data is established in Mn_{1-x}Fe_xSi.

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The extrema of the $\partial \chi / \partial T$ derivative may be used for identification of the magnetic phases with long-range and short-range magnetic order.

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Quantum phase transitions in single-crystal $Mn_{1-x}Fe_xSi$ and $Mn_{1-x}Co_xSi$: Crystal growth, magnetization, ac susceptibility, and specific heat

A. Bauer,¹ A. Neubauer,¹ C. Franz,¹ W. Münzer,¹ M. Garst,^{2,3} and C. Pfleiderer^{1,*}

Classical ferromagnetic equation of state is applicable for Mn_{1-x}Fe_xSi.

Belov-Arrott analysis is possible













There are two quantum critical points, x^* and x_c . The first QC point $x^* \sim 0.11$ corresponds to disappearance of LRO and is an underlying one, which is hidden inside the SRO phase. The second QC point $x_c \sim 0.24$ is a "true" one and marks suppression of the magnetic phase with SRO (chiral spin liquid).

 \Rightarrow Quantum bicriticality in Mn_{1-x}Fe_xSi solid solutions.





Phase with short-range order is formed when fluctuations slow down and freeze.

Fluctuations freeze when their radius reaches critical value controlled by disorder (or defects) in the system.



The model. Calculation details.

Classical fluctuations (CF):
$$R_1 = \frac{a_1}{\left[T/T_c(x) - 1\right]^{\delta}}; \quad \delta \sim 1/2$$

Quantum critical fluctuatio ns (QF): $R_2 = a_2 \frac{T_0}{T}$

Defects and disorder (percolation):

$$R_{s} = R_{c} = \frac{l}{(1 - x/x_{c})^{\nu}}; \nu \sim 0.9, x_{c} \text{ - percolation threshold}$$
$$R_{1,2} = R_{s} \Longrightarrow T_{1,2}(x) \qquad \text{Freezing condition}$$
$$R_{1} = R_{2} \Longrightarrow T_{eq}(x) \qquad \text{Crossover between CF and QF}$$

There is only one fitting parameter $T_2(0)$ (measure of the quantum critical fluctuations strength).

Transition into LRO phase $T_c(x) = T_c(0)(1 - x/x^*)$ Transition into SRO phase : $T(x) = \max \{T_1(x), T_2(x)\}$ $T_1 = T_c(x) \left\{ 1 + \frac{\delta T(0)}{T_c(0)} \left(1 - \frac{x}{x_c}\right)^{\nu/\delta} \right\}$ $T_2 = T_2(0) \left(1 - \frac{x}{x_c}\right)^{\nu}$ $\delta T(0) = T_c(0) \left(\frac{a_1}{l}\right)^{1/\delta}; T_2(0) = T_0 \frac{a_2}{l}$

CF-QF crossover temperatur e for model choice $\delta = 1/2$ $A = \frac{z^2}{z-1} \Rightarrow T_{eq}(x)$ $z = \frac{T_{eq}}{T_c(x)}; \quad A = \frac{T_2(0)^2 T_c(0)}{\Delta T(0) T_c(x)^2}$







The predicted crossover line is observed in the resistivity data due to scattering on magnetic fluctuations (Yosida mechanism).

In the case of static susceptibility these fluctuations are averaged and don't affect $\chi(T)$ data.



Why intermediate phase is formed? Why LRO phase is supressed?

PHYSICAL REVIEW B 75, 064430 (2007)

Microscopic model for spiral ordering along (110) on the MnSi lattice





 $J_1 - nn$ (FM) exchange $J_2 - next nn$ (AFM) exchange $J_3 - third nn$ (AFM) exchange

$$J_2 \cong J_3$$

In Heisenberg paradigm RKKY exchange defines J_1, J_2, J_3 parameters, which may be tuned by variation of the electron concentration.

Frustration try to align spirals along (110).

DM interaction try to align spirals along (111)

Competition between two interactions may lead to loosing of the long-range order and formation of the chiral liquid state.





Yosida mechanism holds in Mn_{1-x}Fe_xSi.

Magnetoresistance maximum evolution with iron concentration.





ESR in Mn_{1-x}Fe_xSi. Results.

T (K)

 $T - T_{SP}$ (K)

60 GHz ESR spectra



(enhancement of spin fluctuations) with iron concentration.

Universal scaling $W(T)/W(T_{SP})=1+a(T-T_{SP})^2$ for all concentrations except quantum critical points x^{*} and x_c.

Violation of the standard Korringa relaxation law $W(T) \sim 1/\chi(T) \sim (T - T_{SP})$.

Weakening of the W(T) temperature dependence just at quantum critical points.





Reasonable fit of the experimental data assuming $T_x \sim T_{SP}$, i.e. $T_x \sim 11$ K for x^* and $T_x \sim 0$ K for x_c .

Even in strong magnetic field QC points x^* and x_c derived in the limit B \rightarrow 0 are still here.

ESR is a right tool to visualize QC points including the underlying QC point.



Magnetic phase diagram of $Mn_{1-x}Fe_xSi$ is driven by the sequence of two quantum phase transitions.

First (underlying) QCP is located at x^* ~0.11 and controlled by exchange effects.

Second QCP is located at $x_c \sim 0.24$ and is disorder controlled. This QC point corresponds to the change of the magnetic system topology at the percolation threshold. For $x > x_c$ ground state is a Griffiths phase.

New line on T-x diagram, which corresponds to a crossover between the classical and quantum fluctuations, is predicted in quantum bicritical model and observed experimentally.

Substitution of Mn by Fe leads to enhancement of spin fluctuations (a route from Heisenberg to itinerant case).

Anomalous spin relaxation is a tool to visualize QC points in ESR experiments.



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Quantum bicriticality in $Mn_{1-x}Fe_xSi$ solid solutions: exchange and percolation effects

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Anomalous spin relaxation and quantum criticality in $Mn_{1-x}Fe_xSi$ solid solutions¹⁾

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